Frictional and elastic energy in gecko adhesive detachment

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Geckos use millions of adhesive setae on their toes to climb vertical surfaces at speeds of over 1 m s⁻¹. Climbing presents a significant challenge for an adhesive since it requires both strong attachment and easy, rapid removal. Conventional pressure-sensitive adhesives are either strong and difficult to remove (e.g. duct tape) or weak and easy to remove (e.g. sticky notes). We discovered that the energy required to detach adhering tokay gecko setae (W_d) is modulated by the angle (θ) of a linear path of detachment. Gecko setae resist detachment when dragged towards the animal during detachment (θ = 30°) requiring W_d = 5.0 ± 0.86 (s.e.) J m⁻² to detach, largely due to frictional losses. This external frictional loss is analogous to viscous internal frictional losses during detachment of pressure-sensitive adhesives. We found that, remarkably, setae possess a built-in release mechanism. Setae acted as springs when loaded in tension during attachment and returned elastic energy when detached along the optimal path (θ = 130°), resulting in W_d = −0.8 ± 0.12 J m⁻². The release of elastic energy from the setal shaft probably causes spontaneous release, suggesting that curved shafts may enable easy detachment in natural, and synthetic, gecko adhesives.

Keywords: gecko; adhesion; energy; friction; biomechanics; tribology

1. INTRODUCTION

Geckos’ climbing ability is nearly unrivalled in nature. Geckos are able to climb microscopically smooth surfaces as fast as other terrestrial animals can run on level ground (Autumn et al. 2006b). Key to geckos’ amazing feats are millions of adhesive setae segmented into scansors on the undersides of the toes. Gecko setae form a smart (Fakley 2001) fibrillar adhesive that is capable of strong attachment and rapid and easy detachment. The work of detachment (W_d; the integral of the forces along the detachment path during unloading) represents an adhesive’s ability to resist failure (Zosel 1985). Highly tenacious pressure-sensitive adhesives (PSAs) such as duct tape are capable of absorbing large detachment energies (W_d) yet also require a similarly large W_d for removal. Ideally, a smart adhesive would be capable of absorbing large detachment energies (W_d) while attached, but require only a small W_d during removal. This study focuses on the apparently contradictory design principles underlying adhesion and removal of gecko setae.

1.1. Conventional pressure-sensitive adhesives

Comparisons between the gecko adhesive and conventional PSAs yield striking similarities and profound differences (Autumn 2006b). Both PSAs and geckos adhere primarily by van der Waals (vdW) forces (Autumn et al. 2002; Creton 2003) and have effective elastic moduli below 100 kPa (Dahlquist 1969; Pocins 2002; Autumn et al. 2006c), but gecko setae stand in stark contrast to PSAs in their anisotropy (Autumn et al. 2006a) and self-cleaning (Hansen & Autumn 2005) properties. PSAs are used commonly for both industrial and home applications. Post-It notes contain a PSA that detaches easily from most surfaces, whereas double-sided tape is used typically for structural applications requiring strong permanent attachment (Creton 2003). However, a single conventional PSA cannot adhere strongly and also detach easily. Yet, the normal operation of the gecko adhesive demands both strong attachment and rapid detachment.

The thermodynamic work of adhesion (γ) for vDW interactions is approximately 50 m J m⁻² (Israelachvili 1992) yet W_d is typically orders of magnitude larger in PSAs (Kinloch 1987; Creton & Fabre 2002) as a result of viscoelastic energy loss occurring from internal friction processes such as cavitation and fibrillation of the adhesive (Creton & Fabre 2002) that lead to material degradation. Gecko setae must maintain their structure and adhesive capability over the course of the animal’s moult cycle, approximately two months (K.A. & N.G. personal observation). The viscous deformation typical of detaching PSAs does not seem to occur in the gecko adhesive, which deforms elastically, not plastically (Autumn et al. 2006c). This suggests that if W_d is indeed large in gecko setae, viscous deformation may not be the primary mechanism of energy loss.
1.2. Do fibrillar adhesives yield large $W_d$?

Measurements of $W_d$ in soft rubbers illustrate that highly elastic systems are capable of dissipating energy upon detachment ($W_d \gg \gamma$; Lake & Thomas 1967). The long-chain polymers that make up soft rubbers are cross-linked at large intervals, thereby eliminating local elastic energy return that would otherwise occur during polymer bond rupture and thus results in large $W_d$ (Lake & Thomas 1967). Jagota & Bennison (2002) noted that the gecko adhesive similarly lacks material cross links and proposed that setal fibrils may dissipate elastic energy, resulting in large $W_d$ (Jagota & Bennison 2002). However, the comparison between rubber polymer chains and gecko setae can only be extended so far, as the elastic energy dissipation mechanism in rubbers limits the minimum energy dissipation that occurs, not the maximum (Lake & Thomas 1967). In soft rubbers, large energy dissipation ($W_d \gg \gamma$) will occur regardless of the detachment method. Thus, this elastic dissipation mechanism would not seem to satisfy the gecko’s smart adhesion requirements by enabling strong attachment, yet hindering efficient detachment. To investigate fully the smart adhesive properties of gecko setae, we must consider their unique anisotropy.

1.3. Anisotropic geometry and function of the tokay gecko adhesive

Tokay gecko (Gekko gecko) setae are approximately 110 $\mu$m in length, 4.2 $\mu$m in diameter, and branch at the tips into hundreds of 200 nm wide spatular pads (figure 1; Ruibal & Ernst 1965; Russell 1975; Williams & Peterson 1982; Autumn et al. 2006c). Setae are curved and oriented at approximately 45°, pointing distally (away from the animal), which produces anisotropic shear and normal forces (Ruibal & Ernst 1965; Russell 1975; Williams & Peterson 1982; Autumn et al. 2006c; Tian et al. 2006).

The geometry and function of the tokay gecko seta suggest that to investigate the adhesive’s detachment mechanics, both friction and adhesion effects must be considered. Seta attachment requires a perpendicular preload followed by a small proximal drag (towards the animal) that loads seta tensily (Autumn et al. 2000). Once setae engage, a proximally oriented shear force maintains setal tension and produces adhesion, possibly by elongating the peel zone of the spatula (Tian et al. 2006). A distally directed shear force releases tension and results in no adhesion. Whole animal adhesion measurements show that the tokay gecko's adhesive detaches when the applied detachment force increases above an angle, $\alpha = 26-30^\circ$ (Autumn et al. 2006a) relative to the substrate. This maximum detachment force angle ($\alpha'$) coincides directly with the maximum setal–substrate shaft angle at which point setae spontaneously detach (Autumn et al. 2006a). Increasing the setal shaft angle or applied detachment force angle above $\alpha = 26-30^\circ$ probably causes fracture of the spatula–substrate bonds (Autumn et al. 2000) and thus detachment. However, it is not known how adhering gecko setae respond to applied displacements at varied detachment path angles ($\theta$).

Climbing geckos stay attached to surfaces because they pull their limbs proximally (towards the animal) during the stance phase of the stride (Autumn et al. 2006b). Proximal shearing of the adhesive generates friction and adhesion maintaining attachment (Autumn et al. 2006a). At the initiation of the swing phase of the stride, geckos relax shear force, detaching the adhesive rapidly in approximately 15 ms with no measured reactant forces (Autumn et al. 2006a,b). The climbing dynamics of the gecko demonstrate the directional (proximal–distal) force control of the gecko adhesive’s stickiness. Strong attachment of the adhesive results from proximal shearing as the gecko pulls inward towards the centre of mass (Autumn et al. 2006a,b). Detachment of the foot probably occurs along a vertical ($\theta = 90^\circ$) or distal ($\theta > 90^\circ$) path, but this remains unknown. Here, we present the first experimental measurement of the amount of energy ($W_d$) required to detach isolated gecko setae over a linear detachment path of angle $\theta$. $W_d$ measured while detaching in the proximal direction ($\theta < 90^\circ$) represents the adhesive’s ability to resist failure while trying to maintain attachment. $W_d$ measured during distal detachment ($\theta > 90^\circ$) represents the ease of detachment.

2. MATERIAL AND METHODS

We harvested and mounted 17 setal arrays from five live non-moulting tokay geckos (G. gecko) using the methods modified from Autumn et al. (2002). Under a
from 30 arrays were detached along a linear path of angles ranging and then dragged proximally to attach. Once adhering, setal
directions. (8 towards animal) and distal (180° as vertical and 180° as distal. Through variation of θ, the amount of proximal (x(θ) > 0) or distal (x(θ) < 0) shear during detachment from a set vertical preload displacement (d) is given by the equation

$$x(\theta) = \frac{d}{\tan(\theta)}.$$  (2.1)

Retraction angles (θ) varied from 30° to 150° in increments of 10° with equal angle tested five times in random order. We chose the range of retraction angles to be symmetric about the vertical axis (90°), and constrained between angles that allowed for investigation of a large range of proximal and distal shear. Adhesion in setal arrays requires a vertical preload followed by a proximal shear to engage setae (Autumn et al. 2006a). A preload is a compressive normal (θ = 270°) displacement to a set depth, while a drag is a proximal shear (θ = 0°) displacement. We calibrated the preload displacement for each array by measuring the position at which point a force change of 5 mN is measured during a 10 μm displacement. This position was then the sample’s zero position and all preload displacements reference from it. In the protocol LDPAvg (LDP, load–drag–pull), we preloaded samples to the average normal ‘working depth’ of a typical setal array, 35 μm, and then dragged 500 μm (N = 10 samples × 5 trials × 13 angles = 650). LDPAvg trials produced adhesive forces prior to detachment typical of previous measurements (Autumn et al. 2006c). Detachment from 35 μm between 30° < θ < 150° resulted in shear displacements between 60.62 μm < x(θ) < −60.62 μm as given by equation (2.1).

We were able to control the amount of adhesion in setal arrays before detaching by increasing preload depths in 1 μm increments from the zero position. We determined the preload depth at which adhesion was maximal by monitoring the relative increase or decrease in adhesion during each incremental preload. In the LDPMax protocol, we preloaded additional samples to the depth that maximized adhesion force in each sample (N = 7 samples × 5 trials × 13 angles = 455). In LDPMax trials, we dragged 100 μm prior to detachment.

Detachment velocity for all tests was 50 μm s⁻¹ along the detachment path until arrays completely separated from the slide, yielding a 1.43 Hz unload cycle. To calculate array stiffness and discard any arrays damaged during the experiments, we conducted purely normal load–unload or load–pull tests at the beginning and end of each set of trials for each sample (Autumn et al. 2006c).

2.2. Detachment measurements

Setal array detachment tests were displacement based. A position-controlled microscope slide engaged mounted samples and then detached along the selected detachment path at a retraction angle θ. Retraction angles are defined using 0° as proximal, 90° as vertical and 180° as distal. Through variation of θ, the amount of proximal (x(θ) > 0) or distal (x(θ) < 0) shear during detachment from a set vertical preload displacement (d) is given by the equation

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2.3. Calculations

A custom spline filtering algorithm in LabVIEW 8 (National Instruments, Austin, TX) compensated for force sensor drift, and a third-order Butterworth FFT algorithm set at 30 Hz in DIAdem (National Instruments, Austin, TX) removed high-frequency noise. To calculate the force and energy requirements during detachment, we evaluated the force–displacement measurements: the negative force–displacement integral of the setal array unloading curve yielded the detachment energy (work required to remove the adhering sample: \( W_d \)) as a function of detachment retraction angle. We calculated the shear elastic coefficients for individual setae by determining a sample’s total shear stiffness during the loading cycle, and then scaling coefficients for individual setae by determining a sample’s total retraction angle. We calculated the shear elastic coefficient for both the setal area (A) and duration (T) of unloading as 

\[
\sigma = \frac{1}{(T \cdot A)} \int F_{\text{total}}(t) \, dt
\]

Results are means ± s.e.m. unless otherwise noted.

3. RESULTS

Isolated tokay gecko (\( G. \) gecko) setal arrays exhibited frictional adhesion forces and normal effective stiffness values consistent with previous setal array measurements (Autumn et al. 2006a,c). Setal arrays averaged 0.93 ± 0.12 mm² in area, and contained approximately 13,392 ± 1728 setae, assuming 14,400 setae mm⁻² (Schleich & Kastle 1986). Setal arrays preloaded to 35 μm depth and dragged (LDPavg, \( N=10 \) trials, 650 trials) had average proximal friction stresses of 111 ± 19.8 kPa (104 ± 30 mN) and adhesion stresses of −30 ± 6.7 kPa (−28 ± 8 mN). Arrays preloaded and dragged at individual maximum adhesive depths (LDPmax, \( N=7 \) arrays, 455 trials) had average friction stresses of 184 ± 23 kPa (171 ± 18 mN) and adhesion stresses of −48 ± 7 kPa (−45 ± 6 mN). Setal array effective elastic modulus (\( E_{\text{eff}} \)) in the normal axis averaged 117 ± 15.6 kPa with an 8 ± 4% decrease between initial and final detachment tests.

3.1. Detachment stress

The largest average detachment stresses (\( \sigma \)) for both loading conditions occurred during 30° detachment (figure 3). Adhering setal arrays exhibited large and relatively consistent \( \sigma \) for \( \theta<110^\circ \), at which point \( \sigma \) began decreasing reaching a minimum at \( \theta=130^\circ \). Maximum \( \sigma \) was 53 ± 7.6 kPa for strongly attached arrays (LDPmax protocol) and 35 ± 9.2 kPa for arrays loaded to a 35 μm preload depth (LDPavg protocol). Minimum \( \sigma \) was 15 ± 1.5 kPa for LDPmax trials and 12 ± 2.5 kPa for LDPavg trials.

3.2. Energy dissipation in the gecko adhesive

\( W_d \) varied similarly as a function of detachment angle (\( \theta \)) regardless of loading method. \( W_d \) was maximal at the lowest retraction angle (\( \theta=30^\circ \)), decreasing as \( \theta \) increased. \( W_d \) was negative above \( \theta=110^\circ \). Maximum \( W_d \) was 5.0 ± 0.86 J m⁻² for maximum adhering arrays (LDPmax) and 3.0 ± 0.81 J m⁻² for average adhering arrays (LDPavg). Negative \( W_d \) signifies that energy was returned upon setal array detachment, instead of being dissipated. Minimum \( W_d \) occurred when detaching the gecko adhesive at \( \theta=130^\circ \) signifying greatest energy return at this angle. Minimum \( W_d \) was −0.8 ± 0.12 J m⁻² for maximum adhering arrays and −0.55 ± 0.09 J m⁻² for average adhering arrays.

3.3. Control of \( W_d \) during engagement

Attachment of setae for both LDPmax and LDPavg testing methods occurred within approximately the first 10 μm of dragging, as setae elongated elastically (figure 4). Further dragging of setal arrays resulted in a change in friction of approximately 10% (figure 4), indicating that the forces present at detachment are largely unaffected by drag length. Setae generated most of the shear stress present at detachment within the first 20 μm (figure 4). Maximally adhering samples had
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4. DISCUSSION

Gecko setae are a unique directional adhesive. Frictionally coupled adhesion forces occur when the gecko pulls its adhesive proximally (towards the animal). Pushing distally causes setae to compress, not adhere (Autumn et al. 2006a). Setae detach when the shaft angle (Autumn et al. 2000), or resultant force vector angle (Autumn et al. 2006a) exceeds $\alpha^* = 30^\circ$. In this study, we varied the ratio of normal to shear displacement as we separated setal arrays from the substrate to yield a linear detachment path of angle $\theta$. We found that $\theta$ determined the amount of energy ($W_d$) dissipated by setae prior to reaching their critical angle of detachment. Our most proximal detachment trajectory ($\theta = 30^\circ$) resulted in the largest energy dissipation during detachment. Adhering setal arrays required approximately two orders of magnitude more $W_d$ to remove than the thermodynamic vdW bond energy $\gamma$. Energy dissipation decreased as the detachment trajectory moved from proximal to distal. At $110^\circ$, $W_d$ became negative, indicating that for certain trajectories energy can be returned when gecko setae detach.

Both conventional PSAs and setae adhere primarily through vdW forces and are capable of repeatable attachment under light pressure. PSAs are soft viscoelastic materials (Pocinus 2002) with Young’s modulus below 100 kPa at 1 Hz (Dahlquist 1969). Setae are composed of rigid beams with material elastic modulus of 1.4 GPa (Peattie et al. 2007) but through fibrillation the gecko adhesive’s effective modulus of 100 kPa (at approx. 1 Hz) is close to that of PSAs (Autumn et al. 2006c). The maximum $W_d$ of gecko setae is comparable to typical energies for PSAs with $W_d$ around 2–3 orders of magnitude larger than $\gamma$ (Kinloch 1987; Newby & Chaudhury 1998). However, even though setae and PSA have similar maximal $W_d$, we found that the mechanisms of energy dissipation in setae and PSA differ.

4.1. Elastic energy

The question of how an elastic fibrillar adhesive may dissipate energy has been approached before. Theoretically, elastic loading of fibrils during detachment can cause energy loss as fibrils eliminate local energy transfer (Jagota & Bennison 2002). Elastic potential is given by $W_{\text{elastic}} = (F^2/2k)$, where $F$ is the pull-off force and $k$ is the spring constant. If $W_{\text{elastic}}$ largely determines $W_d$, then $W_d$ should be proportional to the elastic potential, and therefore the stress: $W_d \propto \sigma^2$.

In contrast, for energy-dissipating detachment paths ($\theta < 110^\circ$), $\sigma$ remained largely consistent while $W_d$ varied greatly, suggesting absence of the stress–energy relationship and that elastic loading during detachment does not occur in adhering setal arrays.

Unloaded setae are curved and have increased axial stiffness when straightened in tension. A seta with length 110 $\mu$m, radius 2.1 $\mu$m and Young’s modulus of 1.4 GPa has a curved axial stiffness of

$$k_{\text{curved}} = C \frac{ER^4}{L} \approx 0.20 \text{ N m}^{-1}.$$  (4.1)
Figure 5. Detachment energy ($W_d$) components for distal angles ($\theta > 90^\circ$) illustrate the elastic energy return during engagement. 

(a) Normal $W_d$ shows that setae adhere ($W_d > 0$) up to $\theta = 120^\circ$, after which energy is returned for greater $\theta$. As the detachment angle becomes increasingly distal, setal tension relaxes earlier along the detachment trajectory allowing for large normal energy return during detachment. 

(b) Shear $W_d$ becomes negative as soon as the displacement becomes distal ($\theta > 90^\circ$), indicating that shaft tension is being relaxed elastically. Detachment at angles above $120^\circ$ relaxes shafts and then begins to compress them axially resulting in an increasing $W_d$.

(for a constant of curvature, $C = 10$; Persson 2003) and a straightened axial stiffness of

$$k_{straight} = \frac{\pi R^2 E}{L} = 176.3 \text{ N m}^{-1}$$  \hspace{1cm} \text{(4.2)}

(Sitti & Fearing 2003). This large increase in axial stiffness $(k_{straight}/k_{setal}) \approx 10^{-3}$ suggests that axial elastic loading of setae can occur only until shafts are pulled straight. Our results suggest that setae store energy elastically only during their attachment phase and not during further loading. This is consistent with prior studies showing that setal adhesion requires a vertical preload followed by a proximal drag to adhere (Autumn et al. 2000). Setae generated most of their frictional adhesion forces within the first $20 \mu$m of dragging (figure 4). Initially, setae loaded elastically with measured shear setal spring constants of $k_{avg} = 0.61 \text{ N m}^{-1}$ and $k_{max} = 1.11 \text{ N m}^{-1}$. The measured shear elastic coefficients are well within the range of axial setal stiffness values predicted from our calculation ($0.20 \text{ N m}^{-1}$, equation (4.1)) and the $6 \text{ N m}^{-1}$ calculation of Persson (2003). Thus, it is reasonable that the stiffness we measured while pulling setae in shear corresponds to the axial stiffness of the curved setal shafts as they extended. Setal forces stabilized after approximately $10 \mu$m of dragging, indicating that as setal stiffness increased from elongation, the tips began sliding and shaft tension equilibrated. For a setal array density of $14 400 \text{ mm}^{-2}$ with a setal stiffness range of $k = 0.61$–$1.11 \text{ N m}^{-1}$, the energy stored during the $10 \mu$m setal array elastic loading was

$$W_{avg} = \frac{1}{2} k_{avg} x^2 = 0.44 \text{ J m}^{-2}$$  \hspace{1cm} \text{(4.3)}

$$W_{max} = \frac{1}{2} k_{max} x^2 = 0.80 \text{ J m}^{-2}$$  \hspace{1cm} \text{(4.4)}

Equations (4.3) and (4.4) illustrate that even if elastic loading of the curved setae occurred upon detachment, it could not account fully for the large values of $W_d$ (up to $5 \text{ J m}^{-2}$). Thus, while elastic loading of the setae is an integral part of the gecko adhesive’s attachment mechanism, it does not seem to cause energy dissipation upon detachment. However, elastic energy does play a central role in energy return for highly distal detachment paths. Equations (4.1) and (4.2) suggest that setae pulled proximally during loading straighten and become effectively inextensible (figure 4), limiting possible elastic energy dissipation. However, elastic unloading of setae did occur during detachment, causing energy to return to the system. The minimum values of $W_d$ represent the maximum amount of energy returned during detachment at path angle $\theta$. Previous tests on polytetrafluoroethylene showed that non-adhering setae return nearly all elastically stored energy during unloading (Autumn et al. 2006c). Setal attachment requires normal and proximal/shear loading, which account for the net elastic energy that is returned when detaching along paths of $\theta > 110^\circ$. Equations (4.3) and (4.4) predict that $0.44 \text{ J m}^{-2}$ ($LDP_{avg}$) and $0.80 \text{ J m}^{-2}$ ($LDP_{max}$) of shear elastic energy was stored during detachment. This is similar to the maximum energy return in the shear direction of $0.56 \pm 0.1 \text{ J m}^{-2}$ ($LDP_{avg}$) and $-0.92 \pm 0.08 \text{ J m}^{-2}$ ($LDP_{max}$) both at $120^\circ$ (figure 5). Vertical loading of the gecko adhesive to a depth of $35 \mu$m with average effective elastic modulus $117 \text{ kPa}$ and approximate height of $70 \mu$m (Autumn et al. 2006c), results in $W_{avg} = (E_{eff} (\Delta H)^2 / 2H_0) = 0.887 \text{ J m}^{-2}$ of normal elastic energy. The maximum energy return in the normal direction for $LDP_{avg}$ was $-0.84 \pm 0.1 \text{ J m}^{-2}$ at $150^\circ$, indicating that approximately 96% of normal elastic energy may be returned at higher angles.

### 4.2. Frictional detachment

Once setae are loaded fully by the load and drag steps, additional shear causes setal tips to slide along the substrate (figure 4). Consistent friction and adhesion forces for drags above approximately $20 \mu$m (figure 4)
show that sliding maintains constant shaft tension. Therefore, detaching setae already in tension causes setal tip sliding and results in frictional energy dissipation. In figure 6, tip sliding is shown for a 30° linear detachment path where substantial sliding between the spatula and substrate occurs as the setae detach. The average sliding friction stress for a setal array was 111 kPa: for just 10 μm of tip sliding during detachment, $W = \sigma \cdot d = 1.1 \text{ J m}^{-2}$ of energy could be dissipated.

Figure 6. Average unloading friction and adhesion curves from setal array detachment tests ($\text{LDP}_{\text{avg}}$) are shown for three detachment angles. For each detachment angle, a sequence of still images from a video of a single seta illustrates the kinematics of detachment at that angle. Positive shear force indicates that setae are pulled in tension, and a negative normal force indicates that setae are adhering. The white line in the single seta image sequences indicates the initial setal tip position prior to detachment, and the grey bar shows the total distance of tip sliding. The last frame of each sequence has a 50 μm scale bar in the upper left corner. Left tip displacement is distal and right tip displacement is proximal. (a) Detaching along a linear path oriented at 30° causes shafts to remain in tension during detachment. The shear force slowly decreases during the largely proximal detachment path until removed from the surface. The setae adhered during the whole detachment motion requiring a large energy input necessary to detach. The single seta images show that the tip slid approximately 30–50 μm prior to detachment. (b) Detachment along a 130° linear path resulted in a rapid decrease of both shear and normal force signifying that shaft tension was released rapidly. The single seta images show no tip displacement during the entire unloading cycle. This suggests that 130° is the optimum detachment angle because it allows elastic shaft energy to return without dissipating any energy frictionally. (c) A linear detachment path oriented at 150° resulted in a rapid decrease of shear force and normal force. However instead of remaining small, the shear and normal forces became substantially large in the opposite direction indicating that the setae was compressing during detachment. The image sequence shows that, at such a high detachment angle, the shaft is compressing axially and sliding during detachment, which resulted in energy dissipation and inefficient detachment.
Viscoelastic energy loss is the dominant mechanism for tenacious adhesion of PSAs. Energy dissipation in PSA detachment occurs as a result of internal friction during the deformation of the PSA medium (Blushan 2002; Creton 2003). In contrast, energy dissipation in the gecko adhesive occurs through external friction at the seta-substrate interface. Thus, the gecko adhesive possesses the ‘viscous’ energy-dissipating properties of PSAs, yet is able to maintain its structure and adhesive capability over many loading/unloading cycles without the material degradation associated with plastic deformation. Frictional energy dissipation has also been shown to occur in the asymmetric detachment of PSA interfaces, where friction accounts for up to 75% of the total fracture energy measured (Newby & Chaudhury 1998). Viscous shear deformation of PSAs is analogous to the ability of setae to absorb large amounts of energy as they slide along the substrate during detachment. However, in addition to dissipating frictional energy during detachment, setal curvature enables efficient spring loaded detachment (figure 7) unattainable in PSAs.

Frictional energy dissipation can be modelled by assuming that a setal shaft in tension acts as an inextensible beam free to rotate. Vertically preloading an approximately 70 μm tall setal array to half its height (d=35 μm; Autumn et al. 2006a) will cause the shaft angle of a $L=110 \mu m$ long seta to rotate to an angle of

$$\alpha_0 = \sin^{-1}\left(\frac{H-d}{L}\right) = 18^\circ.$$  (4.5)

The seta’s initial shaft-substrate angle is $\alpha_0$ and $\alpha$ increases during detachment as the setal base moves along a linear path of angle $\theta$ until $\alpha$ reaches the critical angle of detachment $\alpha^* = 26–30^\circ$ (Autumn et al. 2006a) at which point setae detach spontaneously. The shear tip displacement along a substrate during detachment of an inextensible seta of length $L$ and retraction angle ($\theta$) is

$$s(\theta) = L\left(\cos(\alpha_0) - \cos(\alpha^*) + \frac{\sin(\alpha^*) - \sin(\alpha_0)}{\tan(\theta)}\right).$$ (4.6)

Estimating frictional work as $W_{\text{det}}(\theta) = \sigma_{\text{friction}} \cdot s(\theta)$ where $s(\theta)$ is the tip displacement from equation (4.6), we found that a constant setal friction stress of 161 kPa could account for $W_d$ (figure 8). This stress value is consistent with our drag friction measurement of $111 \pm 19.8$ kPa (Autumn 2006b).

Setae in tension typically slide only proximally because distal base motion results in shaft unloading. Frictional energy dissipation ($W_d>0$) occurs mostly when the tip is sliding proximally during detachment. Using equation (4.6), we calculate that proximal tip displacement of a tokay gecko seta of length 110 μm ($\alpha^*=26–30^\circ$) (Autumn et al. 2006a), preloaded 35 μm ($\alpha^*=18^\circ$, equation (4.5)) should occur only for detachment angles below 112°, at which point tip motion theoretically becomes distal ($s(\theta)<0$). Energy dissipation in the real gecko setal arrays occurred for paths below 110° almost exactly matching this 112° predicted detachment path.

Figure 7. Frictional ($W_d>0$) and spring-loaded ($W_d<0$) detachment mechanisms for an individual seta detaching over a retraction angle $\theta$. Grey arrows illustrate the net energy input (towards the system) and output (away from the system) that occur. (a) During detachment along a 130° linear path, the setal tip remains stationary while the shaft unloads elastically, resulting in spontaneous debonding ($W_d<0$). (b) A largely proximal linear detachment path causes setae to remain in tension, inhibiting elastic energy return to the system. The setal tip slides along the substrate during separation until the setal shaft angle reaches the critical angle of detachment $\alpha^*=26–30^\circ$, at which point debonding occurs. The setal tip sliding that occurs prior to debonding results in frictional energy loss during detachment and thus a large $W_d$.

Figure 8. Frictional tip sliding results in detachment energy dissipation ($W_d$) that varies as a function of $\theta$. Modelling a seta in tension as a rigid rotatable beam $W_d$ can be approximated as $W_{\text{det}}(\theta) = F \cdot s(\theta)$ from the geometric tip displacement $s(\theta)$ of the detaching beam. A seta 110 μm in length is loaded 35 μm from its original height of 70 μm making an initial 18° shaft-substrate angle that increases until detachment occurs at $\alpha^*=26^\circ$. Assuming an average setal stress of 161 kPa produces the fit for $W_d$. The illustration inset in the bottom left corner highlights the frictional sliding of the setal tips during detachment ($\theta<110^\circ$). The illustration inset on the right side shows how negative energy results from the spring-loaded setae unloading during high angle detachment ($\theta>110^\circ$).
4.3. Optimal detachment requires minimizing friction and maximizing elastic return

Setae in their non-adhering default state are curved (Autumn & Hansen 2006), and not yet in tension. Previously, we found that non-adhering setae act as nearly ideal springs (Autumn & Hansen 2006; Autumn et al. 2006). In this study, detaching at increasingly higher angles resulted in relaxation earlier in the detachment trajectory and thus more normal energy is returned (figure 6). However, detaching at paths above $\theta=120^\circ$ also resulted in axial compression, causing an increase in shear energy returned (figures 5 and 6). At $\theta=130^\circ$, there is the best balance of energy return between the shear and normal components. Detaching at $\theta=130^\circ$ causes no tip sliding to occur during detachment and no further energy dissipation occurs (figure 6). Thus, maximum energy is returned during unloading of the setae at the optimal angle of $\theta=130^\circ$.

5. CONCLUSION

Gecko setae adhere strongly to nearly any surface yet detach rapidly and efficiently, unlike conventional PSAs, which are either strong and hard-to-remove or easy-to-remove but weak. For example, duct tape or double-sided tape attach strongly, but large detachment energies ($W_\gamma>\gamma$) are also required to remove these adhesives. Sticky notes detach easily, yet their ability to both adhere strongly (large attachment force) and detach rapidly and efficiently, unlike conventional energy-dissipating systems standard on every automobile. Braking systems rely on traditional frictional mechanics by applying a large normal force to the calliper to create a frictional stopping force. However, a braking system using a synthetic gecko adhesive would not require a large normal load but instead would operate under its own adhesive force, eliminating the need for hydraulic pistons. From our measurements, a gecko brake could dissipate up to $1.8\times10^5 \text{ J m}^{-2}$ for every metre dragged. A 1000 kg automobile moving at 80 km h$^{-1}$ would require $E = (1/2) m v^2 = 2.50 \times 10^6 \text{ J}$ to come to a complete stop. Stopping this automobile in 5 m would require only $A = (2.5 \times 10^5 \text{ J})/((5 \text{ m})(1.8 \times 10^3 \text{ m}^2)) = 0.28 \text{ m}^2$ of gecko adhesive.

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